



Spherical Harmonics-based Visibility Synthesis for Cosmological Intensity Mapping

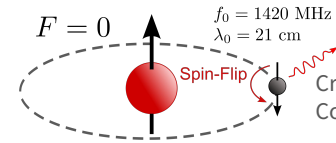
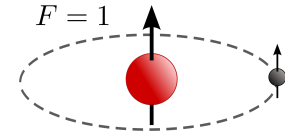
Devin Crichton and Marianna Papadionysiou

ETH Zürich

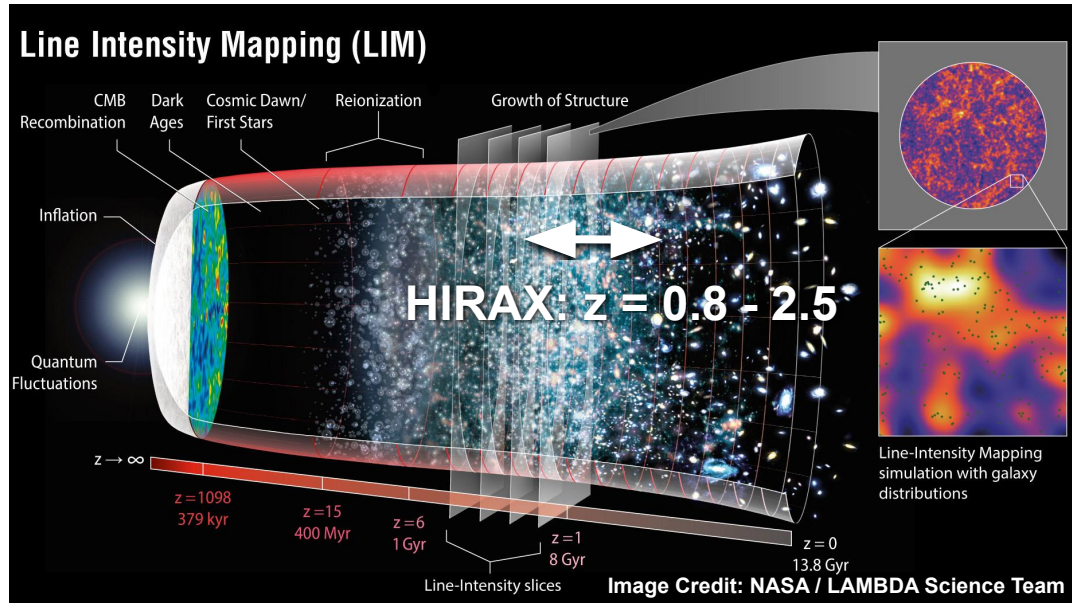
Cosmology in the Alps 2026

HI Intensity Mapping Tomography

- Hyperfine Hydrogen transition line at 1420.4 MHz
- Efficiently and tomographically map cosmological volumes
 - Generally low angular resolution but redshift information cheap
 - Probe epoch of reionisation at low frequencies and large scale structure at high frequencies.



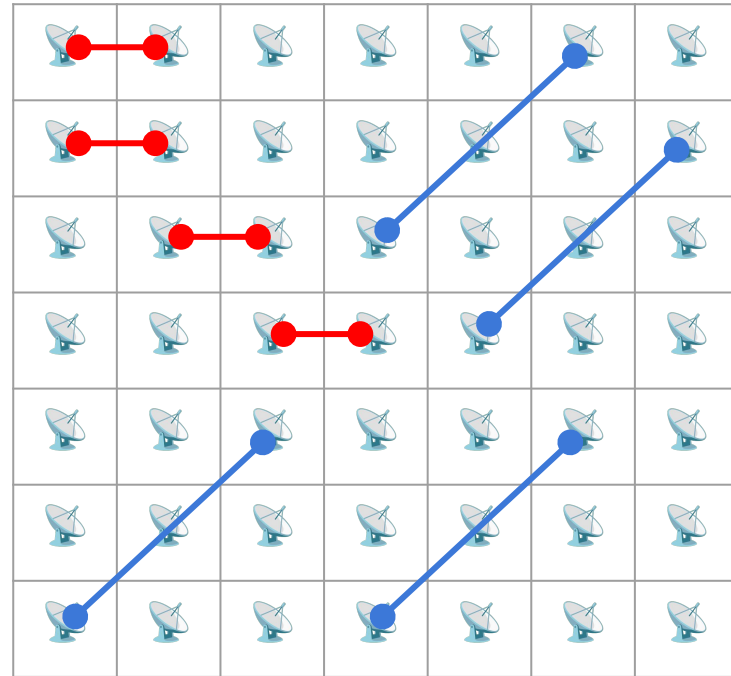
Credit: Wikimedia Commons



Motivation for Compact Redundant Arrays



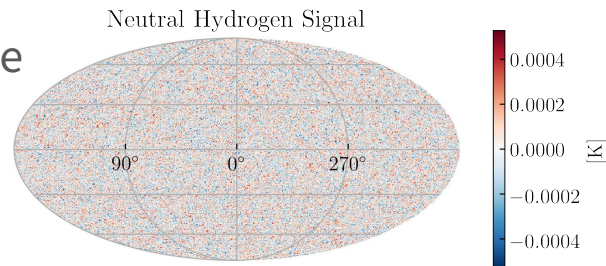
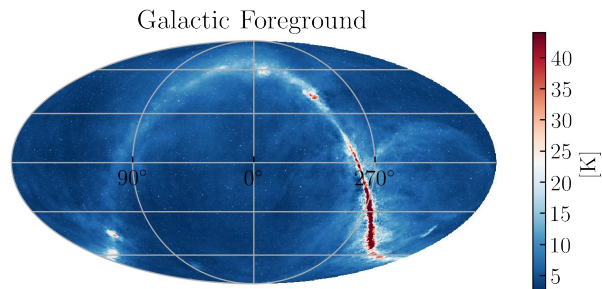
- Compact - weight on short baselines
 - Targeting large angular scales
 - Potential for cross-talk, reflections and impact from array-level effects
- Redundant - Large N with repeated baselines
 - Enhanced sensitivity to sky Fourier modes on interest
 - Internal, redundant calibration
 - If instrument redundancy is realised
 - On-site calibration a requirement for keeping up with data transfer off-site
- E.g HIRAX, CHIME, CHORD, HERA, MWA





Systematics / Chromaticity and Foregrounds

- Foregrounds are the primary challenge for 21cm cosmology
 - Galactic signal brighter by many orders of magnitude
- Signal and Foregrounds have different, *on-sky* properties
 - Galactic emission is:
 - Polarised
 - Strongly correlated over wide frequency bands
 - Structured on the sky in ~known way
 - In principle, there are not many mixed *on-sky* degrees of freedom
- Mode-mixing inherent in measurement/cal. is a major issue
 - Instrument has chromatic response *fundamentally* as well as arising from *systematics*
 - With perfect knowledge of the instrument, this can be accounted for, however the large contrast in signal strengths can make small reconstruction residuals a big problem



Motivation for non-redundant instrument simulation for large -N



Spherical Harmonic Visibility Synthesis

$$\mathcal{V}^{ij\nu}(\phi) = \int d\Omega \overset{\text{Primary Beams}}{A^{i\nu}} \overset{\text{Baseline Term}}{A^{j\nu*}(\hat{\boldsymbol{\theta}})} \exp\left(-2\pi i \mathbf{u}^{i-j} \cdot \hat{\boldsymbol{\theta}}\right) \overset{\text{Sky Signal}}{T^\nu(\hat{\boldsymbol{\theta}}, \phi)}$$

We want to evaluate this integral to forward model for many distinct i and j

Why Spherical harmonics?

- Appropriate for wide field instruments (horizon-to-horizon simulations)
- M-mode separability (see later) / azimuthal sky rotation for **driftscan**
- Analytic tools (e.g. spherical harmonic transform of baseline term)
- Can be compact representation of e.g. beam terms (~azimuthal symmetry)

Standard m-mode (fixed telescope, flexible sky)



Visibility as function
of earth rotation angle ϕ

$$\mathcal{V}^{ij\nu}(\phi) = \int d\Omega \left[A^{i\nu} A^{j\nu*}(\hat{\theta}) \exp\left(-2\pi i \mathbf{u}^{i-j} \cdot \hat{\theta}\right) \right] \left[T^\nu(\hat{\theta}, \phi) \right]$$

SHT [beams x baseline] and sky

$$\mathcal{V}^{ij\nu}(\phi) = \int d\Omega \left[\sum_{\ell_B m_B} B_{\ell_B m_B}^{ij\nu} Y_{\ell_B m_B}^*(\hat{\theta}_{\text{TIRS}}) \right] \left[\sum_{\ell_a m_a} a_{\ell_a m_a}^\nu Y_{\ell_a m_a}(\hat{\theta}_{\text{CIRS}}) \right]$$

$$= \sum_{\ell_B m_B \ell_a m_a} B_{\ell_B m_B}^{ij\nu} a_{\ell_a m_a}^\nu \int d\Omega Y_{\ell_B m_B}^*(\hat{\theta}_{\text{TIRS}}) Y_{\ell_a m_a}(\hat{\theta}_{\text{CIRS}})$$

Azimuthal basis rotation of sky SH's
CIRS \rightarrow TIRS

$$= \sum_{\ell_B m_B \ell_a m_a} B_{\ell_B m_B}^{ij\nu} a_{\ell_a m_a}^\nu \int d\Omega Y_{\ell_B m_B}^*(\hat{\theta}_{\text{TIRS}}) Y_{\ell_a m_a}(\hat{\theta}_{\text{TIRS}}) e^{im_a \phi}$$

Integral of 2 SH's gives delta functions

$$= \sum_{\ell_B m_B \ell_a m_a} B_{\ell_B m_B}^{ij\nu} a_{\ell_a m_a}^\nu \delta_{\ell_B \ell_a} \delta_{m_B m_a} e^{im_a \phi}$$

$$= \sum_{\ell m} B_{\ell m}^{ij\nu} a_{\ell m}^\nu e^{im\phi}$$

TIRS
Fixed to earth
CIRS
Fixed to Sky

Related by pure
azimuthal
sidereal rotation



Standard m-mode (fixed telescope, flexible sky)

$$\mathcal{V}^{ij\nu}(\phi) = \int d\Omega \left[A^{i\nu} A^{j\nu*}(\hat{\theta}) \exp\left(-2\pi i \mathbf{u}^{i-j} \cdot \hat{\theta}\right) \right] \left[T^\nu(\hat{\theta}, \phi) \right]$$

$$= \sum_{\ell m} B_{\ell m}^{ij\nu} a_{\ell m}^\nu e^{im\phi}$$

Fourier transform with respect to earth rotation angle ϕ

$$\mathcal{V}_{m'}^{ij\nu} = \int \frac{d\phi}{2\pi} \sum_{\ell m} B_{\ell m}^{ij\nu} a_{\ell m}^\nu e^{i(m-m')\phi}$$

$$= \sum_{\ell m} B_{\ell m}^{ij\nu} a_{\ell m}^\nu \delta_{mm'}$$

Compact linear map for m-mode visibilities

$$\mathcal{V}_m^{ij\nu} = \sum_{\ell} B_{\ell m}^{ij\nu} a_{\ell m}^\nu$$

M-Modes

Each Fourier component of visibility with earth rotation angle (\sim time) picks out a specific sky m . This is very powerful.

Get full sidereal day of data for \sim same cost of naively computing one time step

Have efficient linear projection from sky to visibilities for full sidereal days.

- **But** to compute the beam transfer matrix, telescope model is fixed in cache
- Flexible telescope (e.g. primary beams) is not efficiently computable

Visibility synthesis for driftscan telescopes



M-mode approach for systematics simulation:

- With known instrument, can efficiently predict data as sky varies
 - But hard to vary instrument without recomputing many harmonic transforms
- Very efficient and m-separability very useful, particularly for inverse problems
- Instrument model baked in to beam transfer matrices and they must be recomputed if it's updated, not ideal for systematics
- Requires primary beams evaluated on same grid as the baseline phase term and sky
 - Can't make use of band limits of individual terms separately.
- Can we does something similar with primary beams flexible?

Triple SH m-mode (as flexible as you want)



$$\mathcal{V}^{ij\nu}(\phi) = \int d\Omega \left[A^{i\nu} A^{j\nu*}(\hat{\theta}) \right] \left[\exp\left(-2\pi i \mathbf{u}^{i-j} \cdot \hat{\theta}\right) \right] \left[T^\nu(\hat{\theta}, \phi) \right]$$

SHT beams, baselines and sky

$$\mathcal{V}^{ij\nu}(\phi) = \int d\Omega \left[\sum_{\ell_p m_p} p_{\ell_p m_p}^{ij\nu} Y_{\ell_p m_p}(\hat{\theta}_{\text{TIRS}}) \right] \left[\sum_{\ell_u m_u} u_{\ell_u m_u}^{i-j, \nu} Y_{\ell_u m_u}(\hat{\theta}_{\text{TIRS}}) \right] \left[\sum_{\ell_a m_a} a_{\ell_a m_a}^\nu Y_{\ell_a m_a}(\hat{\theta}_{\text{CIRS}}) \right]$$

Azimuthal basis rotation of sky SH's
CIRS → TIRS

$$= \sum_{\ell_{pua}, m_{pua}} p_{\ell_p m_p}^{ij\nu} u_{\ell_u m_u}^{i-j, \nu} a_{\ell_a m_a}^\nu \int d\Omega Y_{\ell_p m_p}(\hat{\theta}_{\text{TIRS}}) Y_{\ell_u m_u}(\hat{\theta}_{\text{TIRS}}) Y_{\ell_a m_a}(\hat{\theta}_{\text{TIRS}}) e^{im_a \phi}$$

Integral of 3 SHs gives Gaunt coeff.

$$= \sum_{\ell_{pua}, m_{pua}} p_{\ell_p m_p}^{ij\nu} u_{\ell_u m_u}^{i-j, \nu} a_{\ell_a m_a}^\nu e^{im_a \phi} \mathcal{G}_{m_p m_u m_a}^{\ell_p \ell_u \ell_a}$$

Triple spherical harmonic integral maintains m-mode separability but now we can sum over whatever term(s) we want to fix and keep others flexible

- But need Gaunt coefficients... not as nice as delta functions but not so bad.



Triple SH m-mode (fix sky, flexible beams)

$$\begin{aligned} \mathcal{V}^{ij\nu}(\phi) &= \int d\Omega \left[A^{i\nu} A^{j\nu*}(\hat{\theta}) \right] \left[\exp\left(-2\pi i \mathbf{u}^{i-j} \cdot \hat{\theta}\right) \right] \left[T^\nu(\hat{\theta}, \phi) \right] \\ &= \sum_{\ell_{pua}, m_{pua}} p_{\ell_p m_p}^{ij\nu} u_{\ell_u m_u}^{i-j, \nu} a_{\ell_a m_a}^\nu e^{im_a \phi} \mathcal{G}_{m_p}^{\ell_p \ell_u \ell_a} \end{aligned}$$

Fourier transform with respect to earth rotation angle ϕ

$$\mathcal{V}_m^{ij\nu} = \sum_{\ell_{pua}, m_{pu}} p_{\ell_p m_p}^{ij\nu} u_{\ell_u m_u}^{i-j, \nu} a_{\ell_a m}^\nu \mathcal{G}_{m_p}^{\ell_p \ell_u \ell_a}$$

Sum over baseline harmonics (u) and sky harmonics (a)

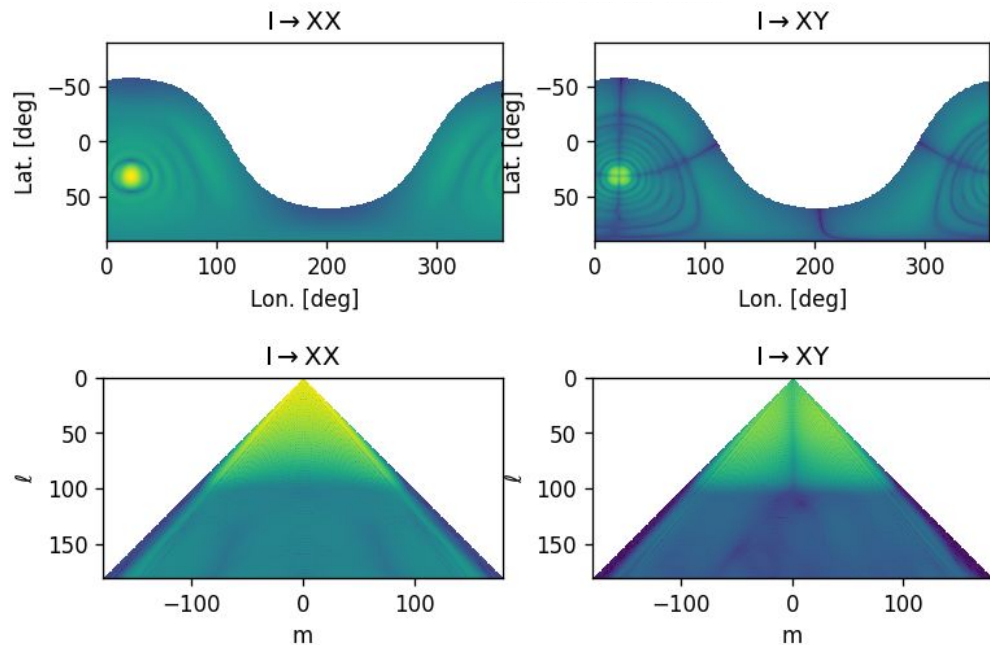
$$\boxed{\mathcal{V}_m^{ij\nu} = \sum_{\ell_p, m_p} P_{\ell_p m_p}^{i-j, \nu} p_{\ell_p m_p}^{ij, \nu}}$$

Here we do what we typically want for systematic non-redundancy studies. Fix the sky and choose some baselines. Then we have an m-separable linear model for the visibilities as a function of power primary beam spherical harmonics.

- Power primary beam harmonics can be generated from voltage beams with similar methods



Aside: Beam m-compactness

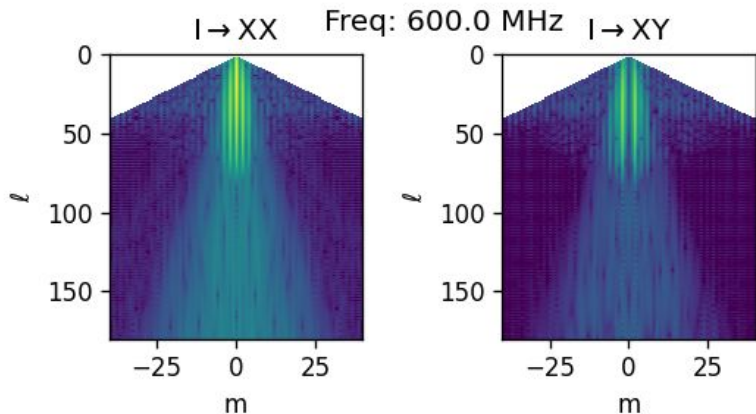
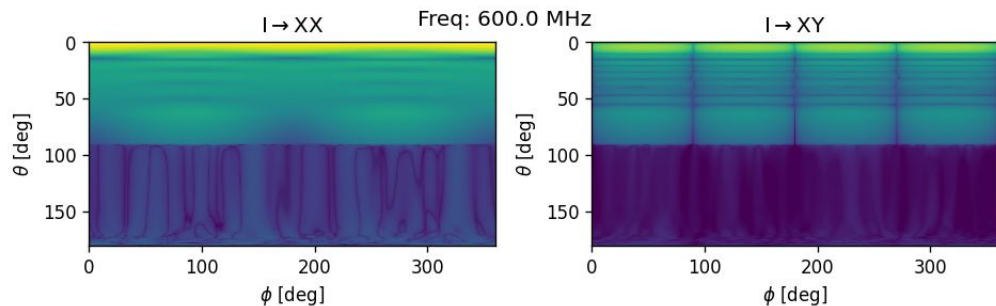


HIRAX beams are compact in l but, in TIRS, they are not compact in m .

However, they are quite azimuthally symmetric around the pointing direction, even e.g. XY terms



Aside: Beam m-compactness



But we can represent them in a pointing-oriented coordinate system.

Then they are compact in l and m . We just need to include a SH rotation term in our cache. (Wigner-D matrices).

If we do so, the cache and beam representations can be extremely compact



SERVAL

Spherical Expansions for Radio Visibility Analysis at Large-N

[CODE](#) | [DOCS](#) - Authors: myself and Marianna Papadionysiou



Python Package with C++ backend

- Computes Gaunt coefficients
- Efficiently sums over them in the ways we want for systematics forward modelling
- Allows Gaunt-based generation of beam transfer matrix or fixed-sky visibility projector
- Misc tools for spherical harmonics and rotations



SERVAL

Spherical Expansions for Radio Visibility Analysis at Large-N

Status: **Early release available. Final release and publication coming soon**

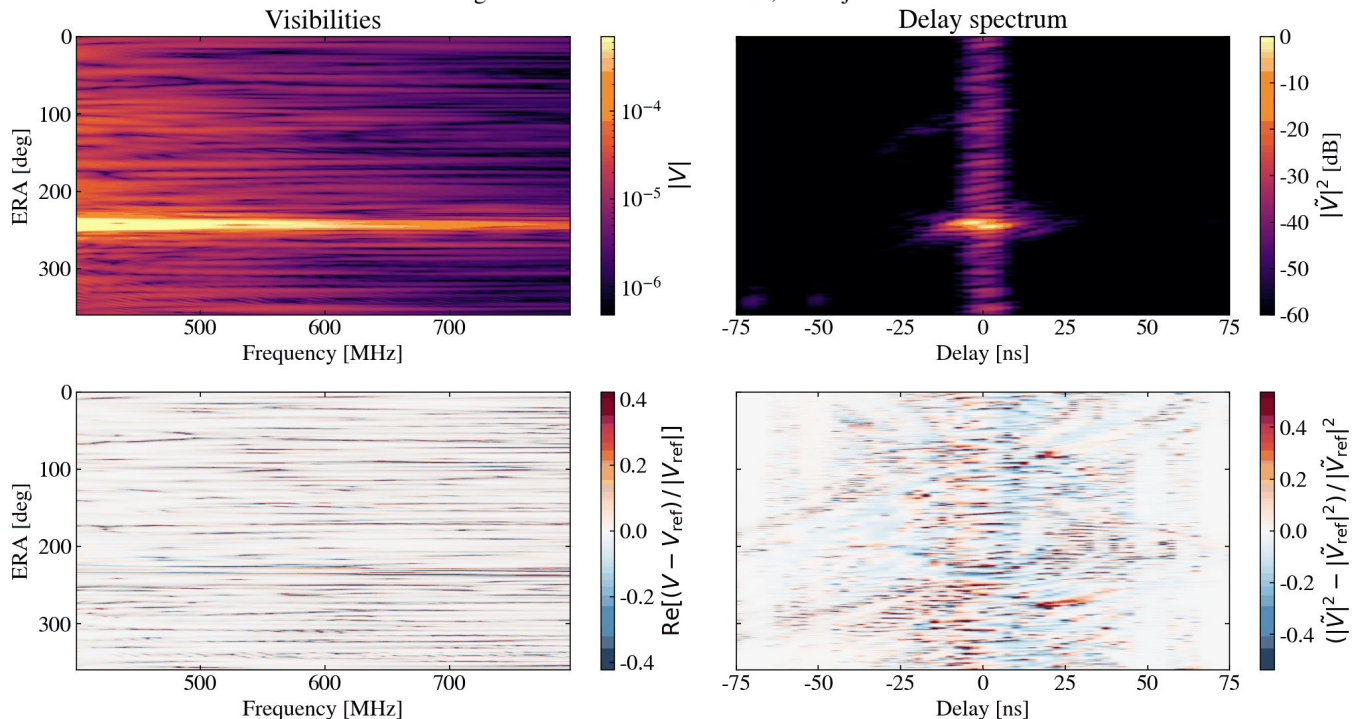
- Mostly feature complete with significant test coverage
 - Includes tests with known analytic results
- Benchmarking and performance improvements ongoing but significant progress has been made in the last months. Performance is very reasonable for typical HIRAX-scale problems
- Documentation is substantial but incomplete and being worked on
- Happy to discuss usage with anyone interested!



Example Results - Pointing Offset Airy Beams

Baseline $E = 13.0$ m, $N = 17.0$ m, $U = 0.0$ m
Pointing offsets: Dish i = 16 arcmin, Dish j = 21 arcmin

Visibilities for a library of voltage beams from pointing-perturbed dishes

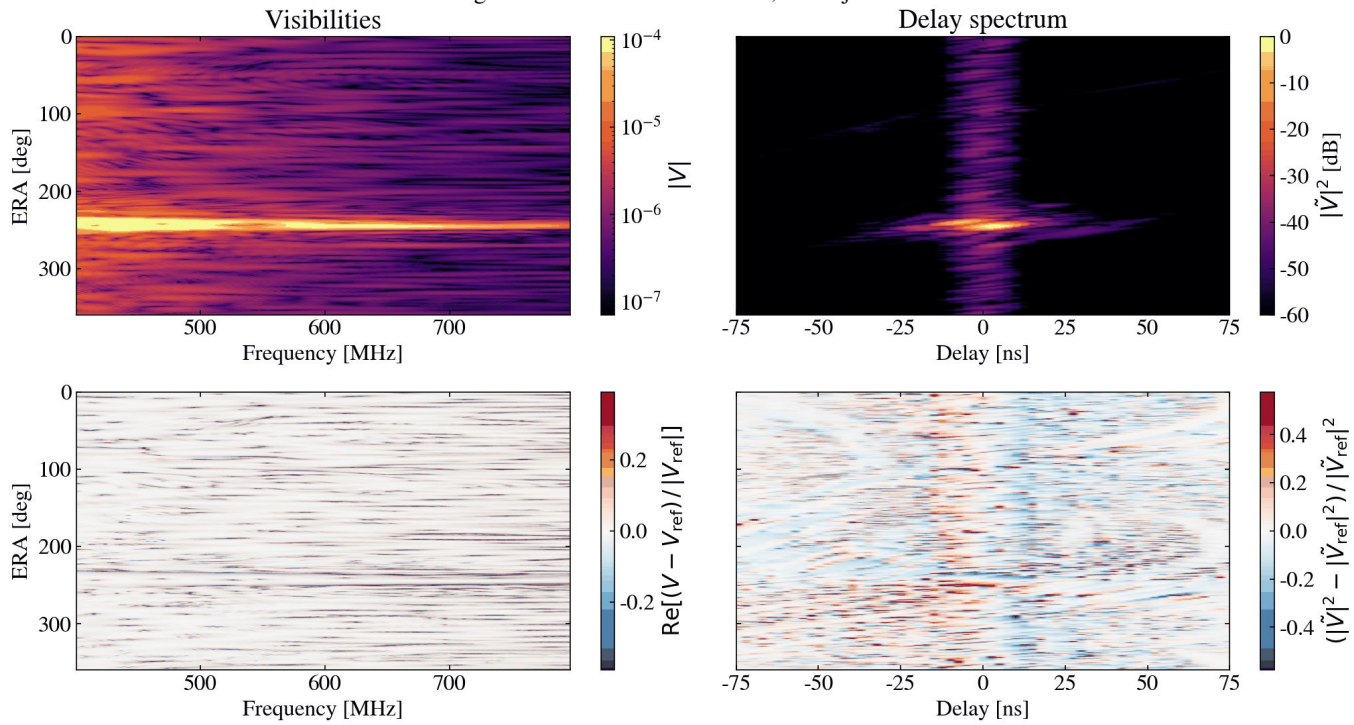




Example Results - Pointing Offset Airy Beams

Baseline E = 26.0 m, N = 34.0 m, U = 0.0 m
Pointing offsets: Dish i = 16 arcmin, Dish j = 21 arcmin

Visibilities for a library of voltage beams from pointing-perturbed dishes





Scaling - Cached vs Direct Visibility Synthesis

Example timings for 128 channel (400-800 MHz) fully sampled sidereal days on 128 core node, power beam lmax: 200

$$\mathcal{V}_m^{ij\nu} = \sum_{\ell_p, m_p} P_{\ell_p, m_p}^{i-j, \nu} p_{\ell_p, m_p}^{ij, \nu}$$

- Visibility generation trivial post cache generation
- Cache generation pays off after a moderate number of beam evaluations (baseline dependent)

Baseline	Direct per beam pair* [min]	Cache Generation [min]	N beam pair direct per cache gen.	N beam pairs HIRAX 128 elem. (full pol.)
E: 6.5 m, N: 8.5 m ($\sim 30\lambda_{\min}$)	3.2	21	7	484
E: 13 m, N: 17 m ($\sim 60\lambda_{\min}$)	4.0	63	16	400
E: 26 m, N: 34 m ($\sim 120\lambda_{\min}$)	6.9	220	33	256



Conclusions

Future and ongoing work:

- Generating a library of caches for HIRAX baselines and standard sky models
- Utilising this to produce visibilities from simulated HIRAX beam responses with systematics perturbations and non-redundancies
- Incorporating into (quasi-)redundant calibration pipelines and e.g. CorrCal model building
- Propagating current understanding of expected HIRAX systematics into visibility-space statistics
- Developing simulations-informed, systematics-hardened foreground mitigation for cosmological analyses.

Thank you!